

Functional renormalization group for few-body physics

Hierarchy in FRG & superselection rule

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Introduction & Motivation

Functional renormalization group (FRG)

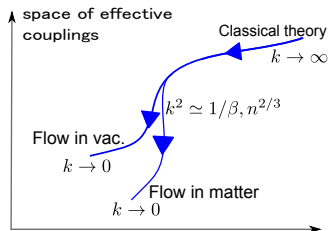
Basic objects:

- Γ_k : low-energy 1PI effective action with a scale $\sim k^2/2m$
- R_k : infrared regulator suppressing low-energy modes

Wetterich equation (Wetterich 1993):

$$\partial_k \Gamma_k[\Phi] = \frac{1}{2} \text{STr} \left[\frac{\partial_k R_k}{\Gamma_k^{(2)}[\Phi] + R_k} \right]$$

($\Gamma_k^{(2)} = \delta^2 \Gamma_k / \delta \Phi \delta \Phi$: field-dependent propagator)



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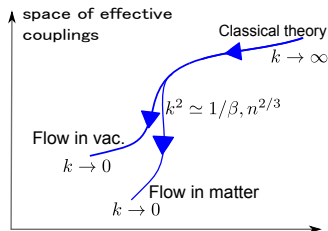
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Features of FRG:

- Non-perturbative approach of quantum field theories
- Practical realization of Wilsonian RG
- **Unified description** of few-body and many-body physics

1-loop property of Wetterich equation

Wetterich equations for self-energy and four-point functions:

$$\begin{aligned}
 \partial_k \rightarrow \square \rightarrow &= \tilde{\partial}_k \left(\text{diagram: square with a loop on top} \right) \\
 \partial_k \left(\text{diagram: square with four external lines} \right) &= \tilde{\partial}_k \left(\text{diagram: two squares connected by a loop} + \text{diagram: two squares connected by a bubble} + \text{diagram: square with a side loop} \right)
 \end{aligned}$$

$(\tilde{\partial}_k = \partial_k R_k \frac{\delta}{\delta R_k}$: k -derivative acting only on explicit k -dependence of regulators)

1-loop property of Wetterich equation

Wetterich equations for self-energy and four-point functions:

$$\partial_k \rightarrow \square \rightarrow = \tilde{\partial}_k \left(\text{one-loop self-energy} \right)$$

$$\partial_k \text{ (four-point vertex)} = \tilde{\partial}_k \left(\text{box with two loops} + \text{box with one loop and two external lines} + \text{box with a loop and two external lines} \right)$$

$(\tilde{\partial}_k = \partial_k R_k \frac{\delta}{\delta R_k}$: k -derivative acting only on explicit k -dependence of regulators)

Important!

- 1PI property + functional realization of RG \Rightarrow **exact** 1-loop property
- Flow of a $2n$ -point vertex depends on $2(n+1)$ -point vertices (**Hierarchy of FRG**)

Schrödinger equation

n -body Schrödinger equation ($|\psi\rangle$: n -body wave function):

$$i\hbar \frac{\partial}{\partial t} |\psi\rangle = (\hat{H}_0 + \hat{V}) |\psi\rangle$$

Hamiltonian manifestly commutes with the number operator \hat{N} :

$$[\hat{H}_0, \hat{N}] = [\hat{V}, \hat{N}] = 0$$

Important!

Particle number operator \hat{N} is a superselection charge.

$\Rightarrow n$ -body scattering amplitude does not depend on $(n+1)$ -body physics

Hierarchy of FRG in few-body physics

Hierarchy of FRG

Flow of a $2n$ -point vertex function depends on $2(n+1)$ -point vertices

However, we must be able to solve this hierarchy for few-body physics **rigorously!**
⇒ We can construct the closed flow equation.

Hierarchy of FRG in few-body physics

Hierarchy of FRG

Flow of a $2n$ -point vertex function depends on $2(n+1)$ -point vertices

However, we must be able to solve this hierarchy for few-body physics **rigorously**!
⇒ We can construct the closed flow equation.

Purpose of this talk

Construction of the closed flow equation for few-body physics.

Solution of the hierarchy
&
Construction of the closed flow equation

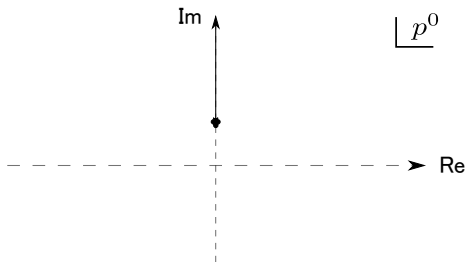
Nonrelativistic model

Propagators (m :mass, μ :chemical potential (< 0)):

$$\longrightarrow = \frac{1}{ip^0 + \mathbf{p}^2/2m - \mu + R_k(\mathbf{p})}$$

Interactions:

$$\begin{array}{c} \diagup \quad \diagdown \\ \diagdown \quad \diagup \end{array} \begin{array}{c} x \\ y \end{array} = \delta(x^0 - y^0) V(\mathbf{x} - \mathbf{y})$$



Two-body physics

Flow equation for four-point vertex functions:

$$\partial_k \left(\text{four-point vertex} \right) = \tilde{\partial}_k \left(\text{diagram 1} + \text{diagram 2} + \text{diagram 3} \right)$$

Each diagram in second and third terms must contain sub-diagrams such as

$$\text{tadpole diagram} = 0.$$

Two-body physics

Flow equation for four-point vertex functions:

$$\partial_k \text{ (four-point vertex) } = \tilde{\partial}_k \left(\text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} \right)$$

Each diagram in second and third terms must contain sub-diagrams such as

$$\text{Diagram with wavy lines} = 0.$$

⇒ Second and **third** terms vanish, and hierarchy is solved for two-body physics:

$$\partial_k \text{ (four-point vertex) } = \tilde{\partial}_k \text{ (Diagram 1) }$$

Three-body physics

Flow of six-point vertex functions:

$$\partial_k \left(\text{square with 6 external lines} \right) = \tilde{\partial}_k \left(\text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} \right)$$

Does the last term vanish as in the case of two-body physics?

Three-body physics

Flow of six-point vertex functions:

$$\partial_k \left(\text{square with 6 external lines} \right) = \tilde{\partial}_k \left(\text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} \right)$$

Does the last term vanish as in the case of two-body physics?

Ans.: No. If we put the last diagram to be zero, we would lose

- consistency with other non-perturbative approach, such as the Dyson-Schwinger (DS) equation,
- universality of renormalization group, such as arbitrariness of R_k , etc.

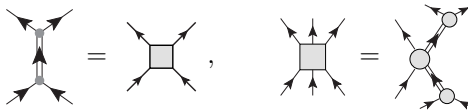
Substructure of eight-point vertex functions

Let us introduce another notation for four-point vertex and following decomposition of six-point vertices:

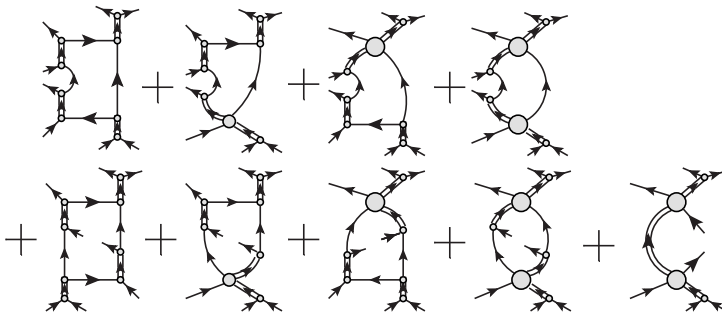


Substructure of eight-point vertex functions

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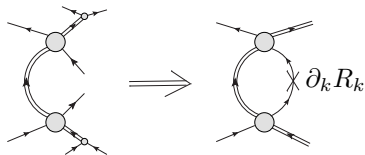
Subset of diagrams in the eight-point vertex functions:



Closed flow equation for three-body scattering problem

$$\partial_k \text{ (diagram)} = \tilde{\partial}_k \left(\text{diagram 1} + \text{diagram 2} + \text{diagram 3} \right) + 9 \text{ feedback terms}$$

Feedbacks are obtained as follows:



Important!

The *closed* and *exact* flow equation for three-body physics is constructed.
 \Rightarrow The solution of this closed flow equation can be shown to satisfy the DS eq.

Summary & Perspectives

Summary

- FRG provides a unified description for few- and many-body physics.
- Hierarchy in FRG can be solved rigorously for few-body physics.
- Construction of the closed flow equation is completed for two- and three-body physics

Perspectives

- General construction of the closed flow equation.
- Application to Efimov physics of the exact and closed flow equation
- Application of knowledge of three-body physics to many-body properties